

# Computers in Physics

## Real-Time Relativity

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# Real-Time Relativity

Walter Gekelman, James Maggs and Lingyu Xu

## An interactive program provides the means of illustrating the effects of high velocity on the appearance of a common object

**T**his article describes a software package which calculates and displays in real time the shape of a cube moving at relativistic velocity in the sunlit world. Examples of its output are presented to illustrate the effects of high velocity on the appearance of a common object.

The cube may be launched from any position and at any angle relative to the observer, but the velocity,  $\beta = v/c$ , is assumed constant. The parameters used in the program may be varied in real time using buttons and knobs. The entire program is menu-driven, and one can choose to view the cube as a Doppler-shifted object or to have each face colored differently to keep track of the large distortions which can occur. This article contains the theory and computational method used to calculate and display the resulting shape. The most important subroutines are contained in the appendix.

### Introduction

In introductory courses on special relativity, students are taught about Lorentz contraction and time dilation with respect to two inertial frames in relative motion. Many students come away with the impression that if an object were moving toward them at an appreciable fraction of the speed of light, it would appear contracted in its direction of motion. It has been understood for many years that this is not the case, and that the Lorentz contraction applies in a world described by measurement with a lattice of clocks and meter sticks.<sup>1</sup> The "observation" of an object in this world rests on analysis of data tapes issued by detectors and clocks within the lattice, long after the object is gone. All parts of an object must be measured at the same time

in order to observe the phenomenon of Lorentz contraction. The difference between a human observer, or a camera, and this type of measurement is that a light sensor, at any given instant of time, detects light which may have originated from the object at very different times. This effect was recognized by several authors<sup>2,3,4</sup> over thirty years ago, and several calculations were done to find the shape of simple objects moving at relativistic speeds, as seen by a human observer.

For several reasons, little of this work has filtered down to the classroom. Students generally have such a difficult time with the concepts of relativity that many instructors feel additional information may lead to an irreversible overload. Also, there has been very little available in visualization tools to dramatically illustrate relativistic effects. A large part of the difficulty that students have with physics is an inability to form a picture which captures the essence of the subject apart from the mathematics in which it is couched. If the problem involves three- (or more) dimensional forms changing in space and time, ordinary blackboard diagrams become nearly useless.

Fortunately, the introduction of powerful graphics workstations is changing this picture. This article describes, in detail, an interactive program which runs and renders, in real time, a cube moving at constant relativistic speeds in any direction with respect to an observer. The code can be straightforwardly modified to deal with any shape or any velocity trajectory. The article is organized as follows. First, we describe the problem in more detail and review some of the ways others have tackled it. Then we present our algorithm for solving it, before proceeding with a description of the program and its user interface and a presentation of several examples of the output. Finally, we list the most important parts of the code.

### The Visual Appearance of an Object Moving at Relativistic Velocities

Consider the emission of light from a simple object, namely a cube, moving rapidly toward an observer in the absence of gravitational fields. As shown in Fig. 1, a

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spherical wave emitted, with the cube at rest, from a point P on the rear surface, is blocked by the back of the cube and is therefore not visible to the observer at point X. Diffractive effects are not considered here. Light rays are normal to the spherical surface and cannot bend around corners. In contrast, when the cube moves rapidly toward the observer at velocity  $v$  along  $x$ , it can outrace most of the expanding light sphere, and the ray emitted from the

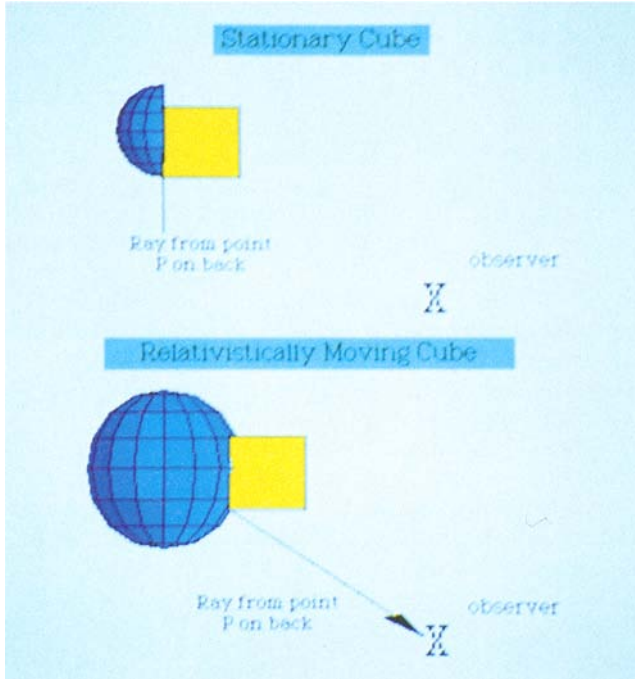


Fig. 1: Illustration of how light emitted from the rear of a rapidly-moving cube can reach an observer. (a) Light emitted from a point on the back of a stationary cube is blocked by the cube. The hemispherical blue surface represents a light pulse emitted from the rear of the cube. (b) A rapidly-moving cube emits a light pulse at the same time and position as the stationary cube in (a), but the cube races to the right so that light from the expanding spherical wavefront is no longer blocked and can arrive at the observer.

point P on the rear face can get to the observer's eye. The angle of elevation,  $\Psi$ , at which this happens is given simply by  $\Psi_\beta = \cos^{-1}(v/c)$ .<sup>5</sup> It is, therefore, possible to see the rear side of a cube approaching at relativistic speed from a viewing angle from which this would not be possible if the cube were at rest!

If the cube is far enough away so that every point on it subtends approximately the same angle  $\Psi$  with respect to the observer, the cube will appear to be rotated as a solid object.<sup>2</sup> The cube appears to rotate as a solid object because it is Lorentz contracted. If it were not Lorentz contracted, it would appear to be elongated along its direction of motion. The relationship between the angle of observation,  $\Psi$ , and the angle of apparent rotation,  $\phi$ , is given by<sup>5</sup>:

$$\phi = \arccos \left\{ \frac{\cos \Psi - \beta}{1 - \beta \cos \Psi} \right\} - \Psi \quad (1)$$

When  $\Psi = 0^\circ$ , the object moves directly toward the observer and the rear face is never observed (i.e.  $\Psi + \phi = 0$ ). When  $\Psi$  increases for a fixed large  $\beta$ ,

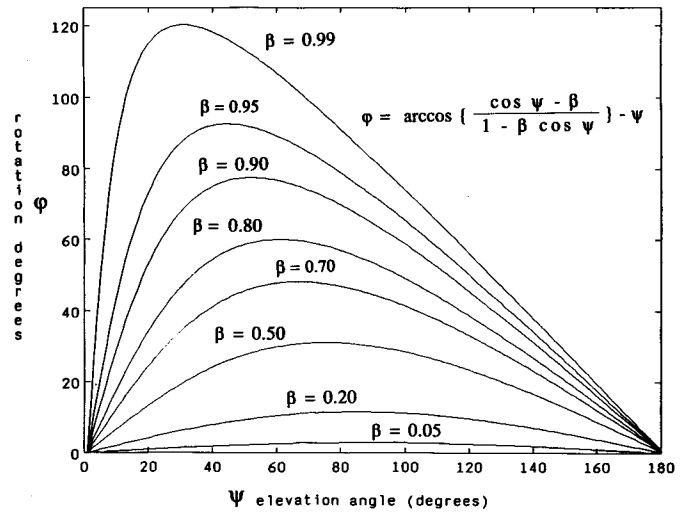


Fig. 2: Graph of Eq. (1), the apparent rotation of a small cube as a function of elevation angle and  $\beta$ . For  $\beta > 0.95$  and a range of elevation angles, the cube can rotate more than  $90^\circ$  so that an observer will see the rear face.

( $\beta = v/c$ ), the rear face of the cube comes into view at the angle of elevation  $\Psi_\beta$  (i.e.  $\Psi_\beta + \phi = \pi/2$ ). In the limiting case as  $\beta \rightarrow 1$ , the cube appears to rotate so that only the rear face is visible from any observation angle (i.e.  $\Psi + \phi = \pi$ ). If one plots the angle of apparent rotation of the cube as a function of the angle of observation<sup>5</sup>, as shown in Fig. 2, another interesting phenomenon emerges. For large  $\beta$  ( $\beta > .95$ ) and a certain range of observation angles ( $\Psi \leq 90^\circ$ ), the cube can appear to rotate more than  $90^\circ$ . In these cases, the bottom face of the cube appears to swap places with the front face, and the rear face with the bottom face, as illustrated in Fig. 3. For small  $\beta$  (0.5), the object rotates slightly ( $\phi = 18^\circ$ ). As  $\beta$  approaches unity, the cube can rotate more than  $90^\circ$  so that the rear is visible to the observer.

The simple analysis breaks down when the cube is close enough to the observer that each point on it subtends a significantly different observation angle  $\Psi$ . In these

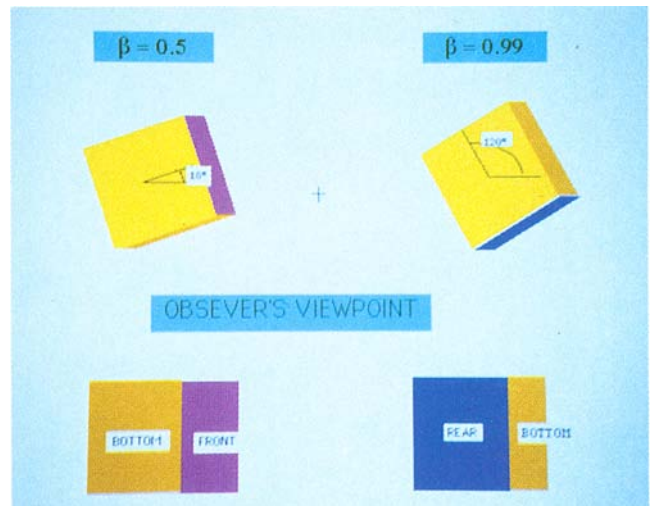


Fig. 3: Illustration of the rotation of a cube which subtends a small solid angle when it moves at intermediate and high  $\beta$ . (a)  $\Psi = 30^\circ$ ,  $\beta = 0.50$ ,  $\phi = 18^\circ$ . (b)  $\Psi = 30^\circ$ ,  $\beta = 0.99$ ; the rear face (colored blue) and the bottom (colored orange) are seen. The rotation angle  $\phi = 120^\circ$ .

circumstances, one could, as suggested by Taylor, approach the problem by breaking the cube up into a multiplicity of smaller cubes and then calculating and performing the above rotation on each cube. This conceptual procedure becomes quite cumbersome in a calculation, since the algorithm must determine how many secondary cubes to break the primary cube into, and then find a way to smoothly join the resulting bunch of differentially rotated secondary cubes for graphical presentation. The problem is further compounded if the original object is a smooth curved surface and not easily represented as a collection of cubes. Because of these difficulties, we choose not to use this approach.

There are other aspects which could be incorporated into a visual presentation. One of these is the Doppler effect, which changes the wavelength of light emitted from the surface of the object according to the relation:

$$\lambda = \lambda_0 \gamma (1 - \beta \cos(\Psi)) \quad (2)$$

where  $\gamma = (1 - \beta^2)^{-1/2}$ ,  $\lambda$  is the observed wavelength, and  $\lambda_0$  is the wavelength emitted when the object is at rest. The cube appears bluer as it directly approaches the observer ( $\Psi = 0$ ), and redder as it recedes ( $\Psi = \pi$ ).

There is also the searchlight effect, in which the distribution of light emitted from a rapidly moving object is most intense along the direction of motion. This effect occurs because, as seen from the observer's viewpoint, the

spherical surfaces containing emitted light energy are closest together along the object's direction of motion, and thus the light intensity is highest in this direction. The distribution of light intensity is given by Weisskopf<sup>2</sup>:

$$I(\theta) = I(\theta') \left\{ \frac{1 - \beta^2}{(1 + \beta \cos(\theta))^2} \right\} \quad (3)$$

with  $I(\theta')$  the angular distribution of light intensity in the object's rest frame. Here,  $\theta$  is the angle of observation of the emitted light with  $\theta = \pi$  being the direction in which the object is moving, and  $\theta'$  is related to  $\theta$  by  $\gamma \sin \theta' = \sin \theta / (1 + \beta \cos \theta)$ . Both the Doppler shift and searchlight effect could be handled using a customized illumination model. However, we do not use this sophisticated approach in the present demonstration because of the difficulty of implementing it with the display software used. The Doppler effect is, though, treated to first order by assuming the entire object is at an average elevation angle, and coloring it uniformly according to the prescription in Eq. (2). In addition, since the color shift is large for rather modest values of  $\beta$ , we arbitrarily limit the amount of color shift to keep the cube in the visible color range.

Finally, there is the effect of relativistic magnification. One sees the rear of a rapidly approaching object in the quasi-remote past, and it appears smaller than the front surface since it was further away when it emitted (or

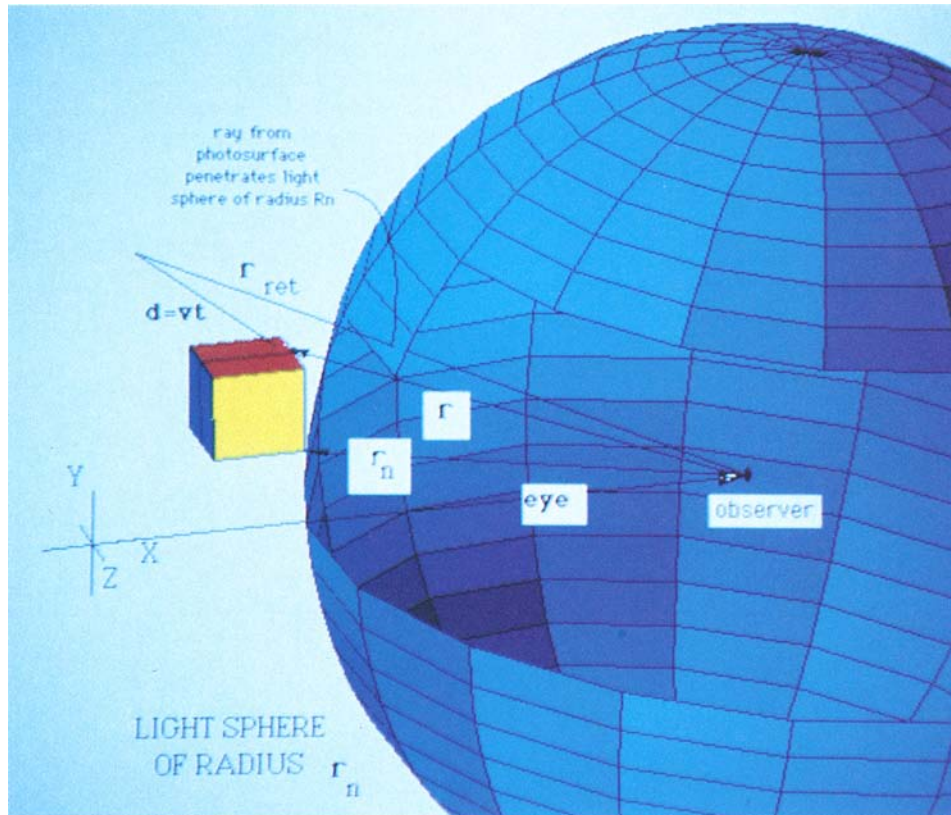


Fig. 4: Some of the vectors used in calculation of point-by-point rotation of surface elements of the cube. The coordinate origin is at  $(x, y, z) = (0, 0, 0)$ , the vector eye goes from the origin to the observer's eye.  $r_n$  is from the observer's eye to the nearest point on the surface of the cube, while  $r$  is a vector to an arbitrary point on the cube.  $r_{ret}$  is the location of the point on the surface (corresponding to  $r$ ) which emitted a light pulse that reaches the observer at the same time as the light from  $r_n$ . The point moves a distance  $d=vt$  in the time it takes the light emitted at  $r_{ret}$  to penetrate a sphere of radius  $r_n$  (drawn in blue) surrounding the observer.

reflected) the light. This magnification effect is present in this demonstration and is noticeable in several of the examples.

The problem is how to handle all these effects in a real-time interactive environment. The ray tracing method circumvents all the calculations involving rotations, by simply following each ray from the eye of the observer back to the object, and keeping track of the different propagation times. This technique has been used by Peterson<sup>6</sup> in an article which contains many striking visual displays. Ray tracing, however, is a time-consuming process which can take from tens of minutes to hours in order to generate one image. There is currently no hardware which can ray trace in real time, so we decided to develop an algorithm, based on ray tracing concepts, which could be implemented on one of the new breed of supergraphics workstations (in this case, a Stardent Titan—64 MB memory, 2 CPUs).

### Calculation of the Appearance of a Relativistically Moving Cube

Light, emitted or reflected from a moving object, reaching an observer's position at time  $t = 0$ , travels various paths of differing length. In order to reach an observer at the same instant, light from a section of the object farthest away from the observer must be emitted earlier than light from the nearest section of the object. Moreover, since the object is in motion, it is in a different position at an earlier time, so that a snapshot of the object in relativistic motion could be distorted from its shape at rest.

Suppose we know the rest shape of an object, and represent it in the computer as an array of points on the surface of the object. The problem then becomes one of computing the spatial location of various points on the object's surface when they emit rays reaching the observer at the same instant, say at  $t = 0$ . To solve this problem, it is conceptually helpful to consider the spatial location of the object at the time the ray traveling the shortest path length reaches the observer. This ray is emitted from the point on the object nearest the observer. The light emitted from the nearest point travels a distance  $r_n$  to reach the observer, in a time interval of length  $r_n/c$ . All other points on the object's surface lie outside a sphere of radius  $r_n$  centered on the observer, as illustrated in Fig. 4. Denoting the position vector of a point on the object's surface, measured from the observer's location, by  $\mathbf{r}(t)$  (vectors are denoted by bold face, with  $\mathbf{v}(t)$  the velocity vector), the object was at position  $\mathbf{r}(t = -r_n/c)$ , moving at velocity  $\mathbf{v}(t = -r_n/c)$ , when the ray from the nearest point was emitted. The trajectory of the cube need not be a straight line moving at constant velocity. For simplicity of analysis, however, we will consider the velocity constant, i.e.  $\mathbf{v}(t) = \mathbf{v}(t = 0)$ . The extension of the method to an accelerated trajectory will be discussed after analyzing the constant velocity case.

Now we calculate the position of some point located on the object's surface (but not the nearest point) when it emits a light pulse that arrives at the observer at  $t = 0$ . The spatial location of this point at time  $t = -r_n/c$  is  $\mathbf{r} = \Delta\mathbf{r} + \mathbf{r}_n$ , where:

$$\Delta\mathbf{r} = \Delta\mathbf{r}_0 - (\Delta\mathbf{r}_0 \cdot \boldsymbol{\beta})(1 - \gamma)/\beta \quad (4)$$

The vector  $\Delta\mathbf{r}_0$  is the displacement vector from the nearest point to the emitting point, measured when the object is at rest. Notice that Eq. (4) contains a Lorentz contraction factor in the direction of the particle velocity, because the location  $\mathbf{r}$  is determined relative to  $\mathbf{r}_n$ , the location of the nearest point, at a fixed time, namely  $t = -r_n/c$ . In order for a pulse of light emitted from the point located at  $\mathbf{r}$  to reach the observer at time  $t = 0$ , it must penetrate the spherical surface about the observer's position of radius  $r_n$  at the time  $t = -r_n/c$ . If the position of the point at the time the pulse is emitted is denoted by  $\mathbf{r}_{\text{ret}}$ , then the time of flight along the ray path, before the pulse penetrates the sphere of radius  $r_n$ , is  $\Delta t = (r_{\text{ret}} - r_n)/c$ , and the emitting point on the object has moved a distance

$$\mathbf{d} = \mathbf{v}\Delta t = \boldsymbol{\beta}(r_{\text{ret}} - r_n) \quad (5)$$

from its location at  $t = -r_n/c$ .

The location of the point at the time of emission,  $\mathbf{r}_{\text{ret}}$ , is related to  $\mathbf{r}$  by:

$$\mathbf{d} = \mathbf{r} - \mathbf{r}_{\text{ret}} \quad (6)$$

where  $\mathbf{d}$  is given by Eq. (5). Taking the vector dot product of both sides of Eq. (6) and using Eq. (5), gives the expression:

$$\beta^2(r_{\text{ret}} - r_n)^2 = r^2 + r_{\text{ret}}^2 - 2\mathbf{r} \cdot \mathbf{r}_{\text{ret}} \quad (7)$$

Taking the vector dot product of Eq. (6) with  $\mathbf{r}$  and using Eq. (5) to replace  $\mathbf{d}$ , the resulting value for  $\mathbf{r} \cdot \mathbf{r}_{\text{ret}}$  used in Eq. (7) gives:

$$r_{\text{ret}}^2(1 - \beta^2) + 2r_{\text{ret}}(\mathbf{r} \cdot \boldsymbol{\beta} + \beta^2 r_n) = r^2 + 2\mathbf{r} \cdot \boldsymbol{\beta} r_n + \beta^2 r_n^2 \quad (8)$$

Equation (8) can be solved for  $r_{\text{ret}}$  using the standard solution for quadratic equations:

$$r_{\text{ret}} = \frac{1}{2a_1} \{ -b_1 + (b_1^2 - 4a_1c_1)^{1/2} \} \quad (9)$$

where:

$$a_1 = 1 - \beta^2 \quad (9a)$$

$$b_1 = 2(\mathbf{r} \cdot \boldsymbol{\beta} + \beta^2 r_n) \quad (9b)$$

$$c_1 = -(r^2 + 2\mathbf{r} \cdot \boldsymbol{\beta} r_n + \beta^2 r_n^2) \quad (9c)$$

The set of all end points of the spatial location vectors,  $\mathbf{r}_{\text{ret}}$ , comprise a surface that we call the *photosurface*. The points on the photosurface correspond to the location of points on the surface of the cube which emit or reflect light that arrives at the observer at the same instant of time ( $t = 0$ , for the case under consideration). The photosurface is generated by the program from the object's location at  $t = -r_n/c$ , using Eqs. (4) through (9). The photosurface can be selected for rendering and viewing from any angle. Of course, the appearance of the object is found only by viewing the photosurface from the observer's location. This view of the photosurface is how the object would appear if a camera located at the observer's position took a snapshot of the object at time  $t = 0$ .

Another way to represent the relativistically moving object is to transform it so that its appearance, when viewed by the observer, is the same as the appearance of the photosurface. We first rotate the vectors (labeled DELTA\_R) leading from the nearest point to each point on the surface of the cube, using the expression for the apparent rotation angle of a cube of negligible size (given by Eq. (1)) located at the retarded position  $r_{ret}$ . The elevation angle of the point located at  $r_{ret}$  is given by:

$$\Psi_{ret} = \arccos (r_{ret} \cdot \beta / \beta r_{ret}) \quad (10)$$

As illustrated in Fig. 5, the rotation of a particular DELTA\_R vector occurs about an axis in the direction of  $r \times \beta$  (which is the same as  $r_{ret} \times \beta$ ). This axis of rotation has been named ROTME. Once the point on the object's surface is rotated in this fashion, the vector from the observer to the rotated point ( $r_{rotated}$  in Fig. 5) is projected onto the direction of the corresponding point on the photosurface. This process ensures that each point on the transformed cube is along the observer's line of sight to the corresponding point on the photosurface. The rotated-projected cube will then appear identical to the photosurface from the observer's viewpoint. The transformed cube is generated by the program using Eqs. (1) and (10), and can be selected for viewing from any angle.

Both the photosurface and rotated-projected cube appear identical only when viewed from the observer's position. However, both objects can be viewed on the computer screen from positions other than the observer's position. Usually their appearance is strikingly different. This ability to view the objects from various aspects can be thought of as having a second observer observing both the original observer and the cube. The second observer's view of the photosurface and transformed cube cannot be realized in the physical world, but does provide some instructive insights into the appearance of the relativistically moving cube.

The procedures used to find the photosurface and transformed cube can easily be generalized to objects of more complex shape and accelerated trajectories. The shape of the object is a problem only in regard to program speed. The data input required is an array of points on the surface of the object at rest. The complexity of the object's shape, or accuracy of its description, is then limited by the array size. Too large an array will slow the program to the point where it can no longer be considered interactive.

An accelerated trajectory can be handled by replacing Eq. (5) with the expression:

$$d = \int_{t_n - t}^{t_n} dt' v(t') \equiv \bar{v}t, \quad (11)$$

with  $t = (r_{ret} - r_n)/c$ , and where  $\bar{v}$  is the average velocity over the interval from  $t_n - t$  to  $t_n$ . The vector  $r_{ret}$  can then be found using an iterative approach. The average velocity is first approximated by setting it equal to  $v(t_n)$ , that is, its instantaneous value when the ray from the nearest point is emitted. The vector  $r_{ret}$  is then found as in the constant velocity case, and the time interval used in Eq. (11) is found using  $t = (r_{ret} - r_n)/c$ . Eq. (11) then gives a new value for  $d$  (and thus  $\bar{v}$ ) which can be used in Eqs. (6) through (9) to find a new value of  $r_{ret}$ ,

and thus a new estimate of the time interval to be used in Eq. (11). This procedure can be repeated until the change in the average velocity after the iteration is below some preset criterion (e.g.  $|\Delta \bar{v}/\bar{v}| < .01$ ).

The velocity trajectory is then broken up into segments, along which the velocity is constant. At one time step, the shapes of the surfaces are computed as described above. At the next time step, in which  $t_n \rightarrow t_n + \Delta t(n)$ , the spatial location of the cube is advanced using velocity  $v_n$ , and the new surfaces are computed as before. Clearly, the repeated procedure for finding the average velocity corresponding to each point on the surface could greatly slow the program. It can be speeded up by replacing the first estimate of the average velocity by the value found at the previous time step for the point in question. In addition, the magnitude of  $\Delta t(n)$  need not be the same for each time step. The value of  $\Delta t(n)$  can be determined, for example, by limiting the size of the derivative of the velocity at each step, i.e. requiring  $|v(t + \Delta t) - v(t)|/|v(t)| < \epsilon$ , where  $\epsilon$  is a small, arbitrarily chosen, positive number. In this case, the value of  $\Delta t(n)$  varies for each time step, and can adequately represent the motion when the acceleration is large.

### Running the Reatiview Program

The user interface for the relativity program is structured so that only a mouse and a dial box are used. Once the program is initiated, by typing RUNME from the control console, a main window and several border windows appear, as shown in Fig. 6. The biggest window, which is positioned in the upper left-hand part of the screen, is the DORE window. DORE (Dynamic Object Rendering Environment) is an object-oriented software graphics system (a product of the Stardent Computer Company, Sunnyvale, CA. It is written in C and portable to other UNIX computers). All the objects required, such as the cube and the "gun" which fires it, are rendered within the DORE window. The cube at rest is shown with the bottom face colored white and the face nearest the observer colored magenta. A three-dimensional grid centered at  $(x = 0, y = 0, z = 0)$  is displayed. The rectangular

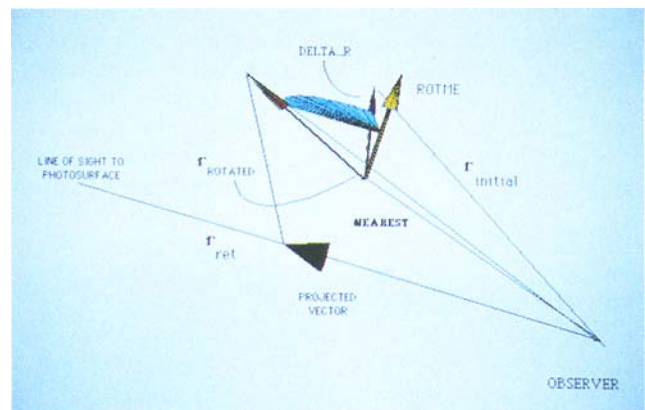


Fig. 5: Elements in the procedure used to transform the rotated cube so that it appears, to the observer, identical to the photosurface. A vector on the surface  $r_{initial}$  is rotated about the axis  $rotme$  by an angle prescribed by Eq. (1). The rotated vector,  $r_{rotated}$ , is then projected onto  $r_{ret}$ , the line of sight to the corresponding point on the photosurface.  $\Delta r$  is the difference between  $r_{initial}$  and nearest, the vector from the observer to the nearest point on the cube ( $r_n$ ).

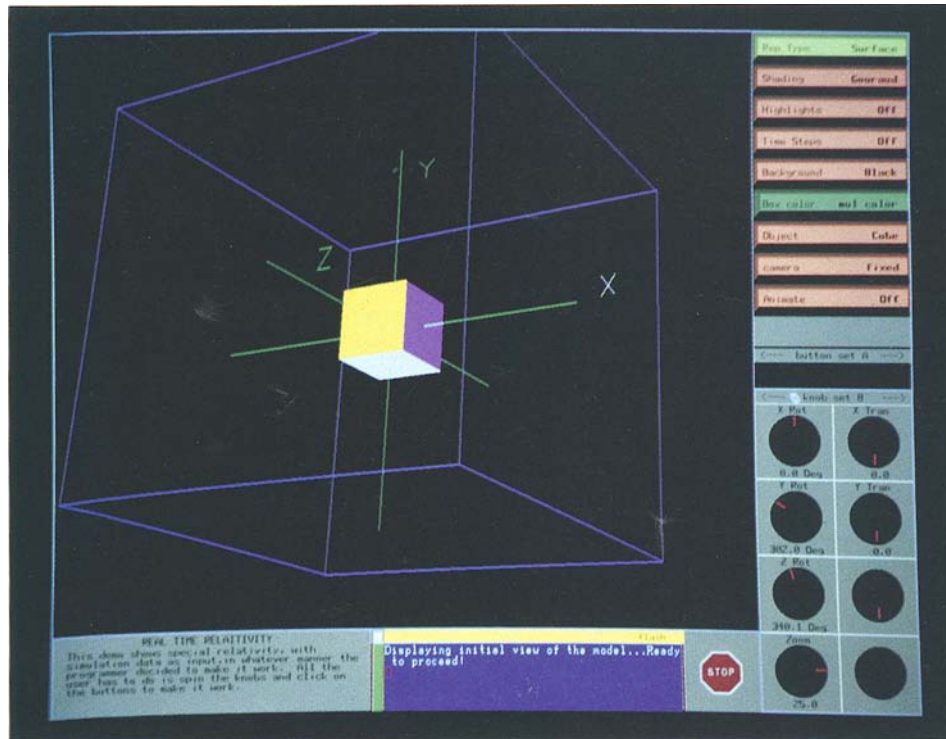


Fig. 6: The computer terminal, showing the screen layout used in the program. A cube is shown at rest at the origin from the perspective of the second observer.

coordinates range from  $-100 \leq x, y, z, \leq 100$ . The unit of distance in the demonstration depends upon the value the user assigns to each tick of the clock. If the time interval is one second, the unit of length is one light-second (the distance light travels in a second), or  $3 \times 10^8$  m. If the time interval is one nanosecond ( $10^{-9}$  seconds), the unit of length is 30 cm.

In the bottom left-hand part of the screen is an explanation/instruction window containing a brief abstract of the program. To the right is an I/O (input-output) window which is initially blank. It has been used during program development for debugging purposes and can show the instantaneous value of a parameter of interest, such as  $\Psi$ , the average elevation angle in real time. A stop sign, displayed in the lower right-hand part of the screen, is used to exit the program. To exit, the mouse is positioned within it and clicked. A button window is located in the upper right-hand part of the screen. To "press" a button, the mouse arrow is positioned on one and clicked. The button functions are explained in detail below. A window in the lower right-hand part of the screen contains a set of dial icons. There is a one-to-one correspondence between the dial icons and the physical dial set. The function of each dial is written on the screen above it. If dial hardware is not present, a dial may be activated by placing the mouse on the icon and clicking. Since there are more dial functions necessary to run the program than physical dials, there are two sets, A and B. Clicking on the button labeled "knob set B" toggles between the two dial sets.

### Buttons

The buttons displayed on the right-hand side of the screen (Fig. 6) perform the following functions:

- (1) Rep Type: sets the mode in which solid objects are drawn. Objects can be displayed as a collection of points, a wireframe structure, or a shaded surface.
- (2) Shading: sets the shading type. The cube and gun can be flat or Gourard-shaded.
- (3) Highlights: determines whether glossy highlights will be present.
- (4) Time steps: sets the maximum number of time steps for the animation. This can be any positive integer number. The default number is 100. If an object moves slowly, it may not go far in 100 steps.
- (5) Background: sets the screen background color to black, red, green or blue.
- (6) Box color: sets the color of the cube to either a single color or a separate color for each face. The single color, which is set to green when  $\beta = 0$ , is used to illustrate the Doppler effect. The separately-colored surfaces are not Doppler-shifted. They are, however, useful when studying the rotations and extreme distortions of the cube at large  $\beta$  and arbitrary direction of motion with respect to the observer.
- (7) Object: determines the object displayed on the screen. In one instance it is the photosurface, and in the second it is the rotated-deformed cube.
- (8) Camera: switches the viewpoint between two positions. In one case the camera is positioned at the observer's location. In the second case the camera may be positioned anywhere, and the observer's position is denoted by an X along the x-axis. This case corresponds to an observer able to observe the original observer.
- (9) Animate: starts or stops a clock determining the time intervals between successive positions of the cube. When the cube is moving and the animate button is pressed, the cube will remain frozen at its last position.

## Dial Set A

All the dials in this set are used to set the cube's initial position and velocity in three-dimensional space. A gun (or launcher) is shown with its muzzle pointing in the direction of the cube velocity.

Dials 1, 2 and 3: set the cube's initial  $x$ ,  $y$  and  $z$  position on the rectangular coordinate axis.

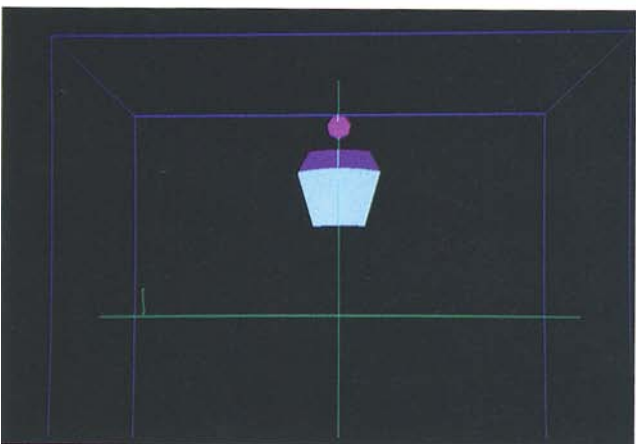
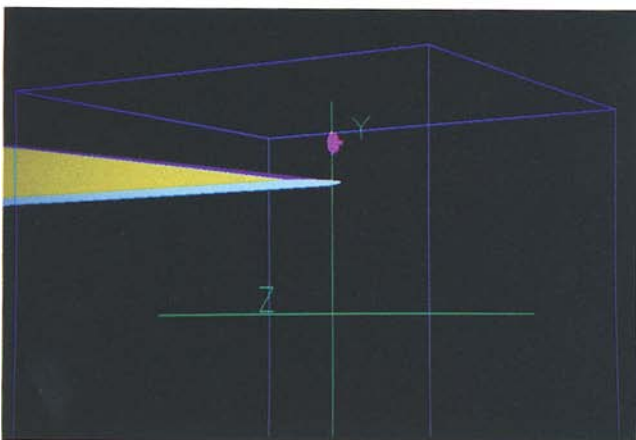
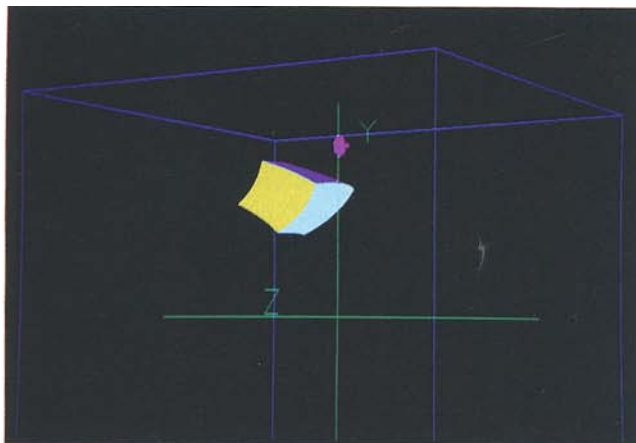


Fig. 7: (a) top; (b) center; (c) bottom. A view of the transformed cube and photosurface from the perspective of the second observer, and also from the observer's viewpoint. The center of the object is at  $y=80.0$ , the viewer is at  $x=600.0$ , and the cube is moving along the  $x$ -axis with  $\beta=0.99$ . This and subsequent figures are discussed in more detail in the "Examples" section of the text.

Dial 4: sets the observer's position on the  $x$ -axis. The current observer's location is displayed by the symbol  $X$ .  
Dial 5 and 6: set the spherical coordinate angles  $\theta$  and  $\phi$ . The angles are in degrees and appear below the knob as it turns.

Dial 7: This dial sets the magnitude of the cube's velocity ( $0.0 \leq \beta \leq 0.999$ ).

Dial 8: This dial changes the intensity of the lights which illuminate the cube.

## Dial Set B

All the dials in this set are used to control the viewing orientation of the second observer. They do not change any of the parameters of the cube motion.

Dials 1, 2 and 3: rotate the entire grid around the  $x$ ,  $y$  or  $z$  coordinate axis. The degrees rotated are shown below the corresponding icon and may take on positive or negative values. Positive rotations are determined by the right-hand rule.

Dial 4: zooms the camera in or out from the grid origin.

Dials 5 and 6: translate the entire grid relative to the center of the screen in the  $x$  and  $y$  directions.

## Examples

Figure 7a shows the appearance of a cube traveling at  $\beta = .99$  (i.e. 99% of the speed of light) in the  $x$  direction toward an observer located at  $x = 600$ . The cube is initially located at  $(x, y, z) = (0, 80, 0)$  so that its elevation angle is  $7.5^\circ$ . Fig. 7b shows the photosurface for the same case as Fig. 7a, from the perspective of an observer looking at both the cube and the first observer.

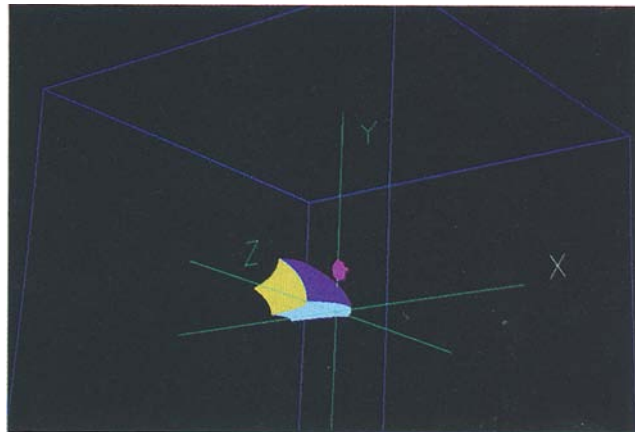


Fig. 8: A cube moving along the  $x$ -axis with  $\beta=0.95$  and center at  $20.0$  shows the differential rotation due to the changing viewing angle across the surface of a large object. The observer is at  $x=100.0$ .

Note that the photosurface is highly elongated, since  $\beta$  is large, but, as illustrated in Fig. 7c, both the cube and photosurface appear identical from the perspective of the observer located at  $x = 600$ . Even though  $\beta$  is large, the apparent rotation (Fig. 3) is less than  $90^\circ$ , since the angle of elevation is small, and the observer sees the (white) bottom and (magenta) front surfaces of the cube.

Fig. 8 shows the distortion of a large cube due to variation of elevation angle from bottom to top. The cube center is at  $y = 20$ , while its lower edge is at  $y = 2$ , and the observer is much closer, at  $x = 100$ . The top edge of the



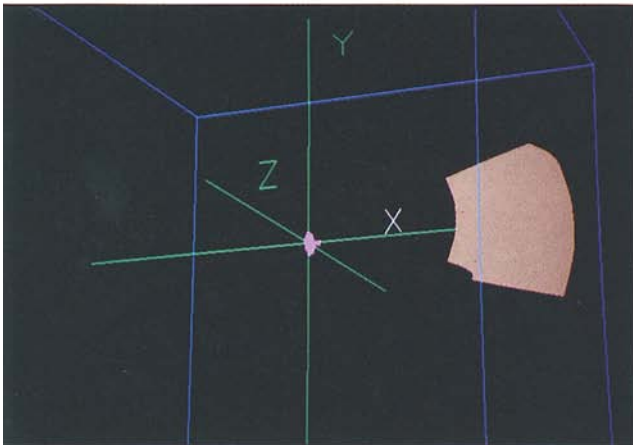
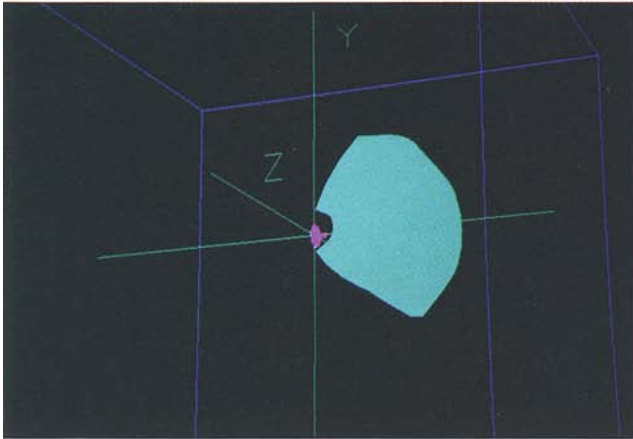
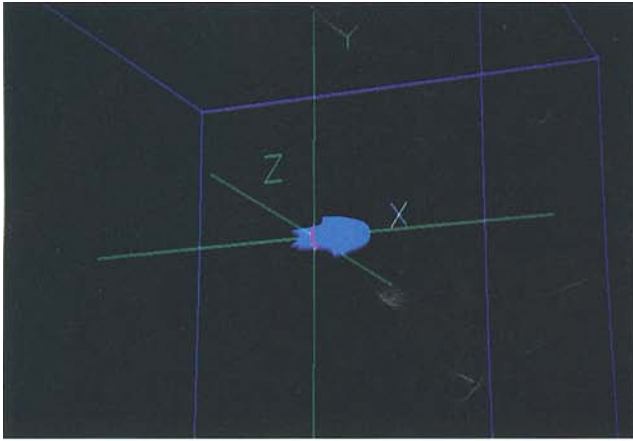


Fig. 9: (a) top; (b) center; (c) bottom. A sequence showing the cube moving along the x-axis through the observer's position at  $x=20.0$  with  $\beta=0.95$ , illustrating both the dramatic change in the appearance of the object and the Doppler shift.

cube corresponds to an elevation angle of  $11.3^\circ$ , and with  $\beta = 0.950$ , a small cube would appear to rotate by about  $55^\circ$ . The bottom edge corresponds to an elevation angle of only  $1.1^\circ$ , so a small cube would rotate about  $5^\circ$ . The large cube in this figure rotates differentially, and thus its front surface becomes rounded.

In the next sequence, a cube, centered at the origin and moving along the x-axis with  $\beta = 0.95$ , approaches an observer at  $x = 20$ . Before the cube reaches the observer (Fig. 9a), it is blue-shifted. At  $\beta = .95$ , the color shift is

enough to move the color out of the visual range. However, the program arbitrarily limits the color shift to keep it in the visual range. The cube appears elongated and nearly bullet-shaped because light from the back of the cube reaches the observer from the past, when the cube was much farther away, so that the rear is squashed down. This is a manifestation of the relativistic magnification effect mentioned earlier. In Fig. 9b, the cube is on top of the observer, and its overall color shifts toward green, which is the object's color at rest ( $\beta = 0$ ). The leading edge appears large, since it is beyond, but close to, the ob-

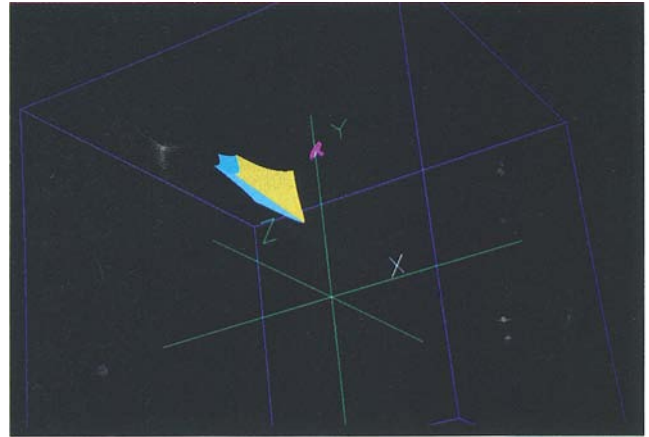


Fig. 10: The cube can appear highly distorted for certain viewing aspects and high velocities, as illustrated here by a cube with velocity direction angles  $\theta=112.5^\circ$ ,  $\phi=43.5^\circ$  and  $\beta=.95$ .

server, so that it subtends a large solid angle. Finally, the cube turns red and appears more cube-like (Fig. 9c) as it recedes from the observer.

Finally, for the same  $\beta$  and observer position as in the previous case, the cube can assume very distorted appearances when it is launched at an arbitrary angle, such as  $\theta = 112.5^\circ$  and  $\phi = 43.5^\circ$ , as shown in Fig. 10. The velocity direction angles are the standard angles used in spherical coordinates, with  $\theta$  measured from the positive z-axis and  $\phi$  measured from the positive x-axis.

### Conclusions

The primary visual effect of relativistic motion viewed in the sunlit world is an apparent rotation. An object moving with velocity  $\mathbf{v}$  and viewed at position  $\mathbf{r}$  appears to rotate about an axis in the  $\mathbf{v} \times \mathbf{r}$  direction. The amount of rotation depends upon the viewing angle and relative speed. In addition, objects which subtend a substantial solid angle with respect to the observer may appear to be distorted. This distortion results from a differential apparent rotation and is not due to a fundamental physical effect such as Lorentz contraction.

The appearance of the relativistically moving object can be modeled by different surfaces. The photosurface is not a rotated object and involves displacements only in the direction of the velocity. The transformed cube involves both rotations and projections. Both objects appear identical from the observer's viewpoint, and are therefore interchangeable representations in this respect. However, to a second observer (who is only metaphysically realizable and exists only in the computer world), the



```

cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
c   Input parameters:
c   eye(3)      position of the observer's eye
c   orig(3)     position of the origin
c   v(3)        velocity of cube  0≤v≤.995
c   pos_gun(3)  position of gun which fires the cube
c   vvv(3)      θ, φ, v  angles gun makes and mag of velocity
c   it          time step ≥ 0
c   Output parameters:
c   surface     surface of distorted cube
c   ph_surface  photosurface of cube
c   hue         Doppler shifted color of cube
c
c   The horizon is at y = 0
c   Written W. Gekelman, J. Maggs, L. Xu 1989
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
c
c   Declaration of variables:
c
c   real eye(3),orig(3)      ! students eye position
c                           ! and origin of coorindate system position
c   real pos_gun(3)         ! init position of cubes
c                           ! center
c   real eye_mag            ! magnitude of eye
c   real v(3)               ! cube's center velocity
c   real vvv(3)             ! theta , phi , v_mag
c   real t,temp             ! time
c   real theta              ! angle as seen by observer
c   real del_theta         ! angle of rotation of each
c   real beta,c             ! v/c
c   real theta_prime,time  ! angle seen by observer
c   real pos_old(3),pos_new(3) ! original position of cube
c   real hue,hue0          ! color 0.0-0.5
c   real shift              ! normalization for color
c   real angle              ! average angle cube makes with horizon
c   real gamma              ! relativistic factor
c   real r_ret_hat(6,16,3) ! unit vector in r_ret direc.
c   real surface(6,16,3)   ! points on object's surface
c   real ph_surface(6,16,3) ! photosurface
c   real surface_orig(6,16,3) ! created surface
c   real Lorentz_surf_orig(6,16,3) ! surface with contractions
c   real Lorentz_surf(6,16,3) ! surface with contractions
c   real r_surf(6,16,4)    ! store scratch rpoint,magr
c   real rotme(6,16,3)    ! rotation axis
c   real sigh_tot(6,16)   ! total rot angle each point
c   real nearest(4)       ! nearest point to eye
c   integer it,i          ! time step
c   save surface_orig,Lorentz_surf_orig ! keep upon reentering
c
c   *****
c
c   The original position of the cube is stored in the
c   1st of 16 arrays
c
c   hue0 = 0.25          ! 0.25=green for color map
c   c = 1.0              ! normalized speed of light
c   if (it.eq.0)then     ! is it the first pass?
c     beta = 0.0
c     do i = 1,3
c       ! save original position
c       ! of the gun
c       pos_old(i) = pos_gun(i) ! initialize position
c       pos_new(i) = pos_gun(i) ! ditto
c       beta = v(i)**2 + beta
c     enddo
c     beta = sqrt(beta)/c
c     if(beta.ge.1.0)beta = 0.99999 ! protection on divide
c     gamma = sqrt(1.0-beta*beta) ! the standard
c     gamma = 1.0/gamma          ! definition of relativistic factor
c
c   Calculate position of points on surface given pos_new and beta
c   First create the cube with center at the orgin
c
c     call create(surface,Lorentz_surf,beta,gamma) ! create cube
c
c   Rotate cube so that face 6 is in beta direction, and place it at
c   initial position pos_new position of cube center
c
c     call rotate(Lorentz_surf,v,beta) !Lorentz surface
c     call rotate(surface,v,beta) ! rotate uncontracted cube
c
c   Next: point cube in beta direc. and translate to ipos()
c
do isur = 1,6          ! save newly created surfaces
do ipnt = 1,16        ! 16 points/side of cube
do ix = 1,3           ! x , y z
  surface_orig(isur,ipnt,ix) = surface(isur,ipnt,ix)
  Lorentz_surf_orig(isur,ipnt,ix) = Lorentz_surf(isur,ipnt,ix)
enddo
enddo
enddo

endif                ! it = 0 surfaces done

time = it             ! make it real
do i = 1,3
  pos_new(i) = pos_old(i) + v(i)*time
  ! for visible motion on screen
  pos_gun(i) = pos_new(i) ! pass position of cube center
enddo

c
c   Move surface to next position Dist = vel*time
c
do isur = 1,6        ! new position of undistorted but Contracted
  ! surface
  do ipnt = 1,16
  do ix = 1,3
    Lorentz_surf(isur,ipnt,ix) = Lorentz_surf_orig(isur,ipnt,ix)
    1 + pos_new(ix) ! move em out
    surface(isur,ipnt,ix) = surface_orig(isur,ipnt,ix) ! rawhide!
  enddo
  enddo
enddo

c
c   Find the position vectors of points on the surface relative
c   to the observer - store them in r_surf.
c
call calc_r(Lorentz_surf,r_surf,beta,eye)

c
c   Calculate axis of rotation for each point on the surface.
c
call axis_of_rot(r_surf,v,beta,rotme)

c
c   Find nearest point to observers's eye and keeps track of it
c   nearest is the point on cube through which overall rotation will
c   be done, ie rotation will be done about axis in rotme direction
c   and point of rotation at nearest!
c
call nearest_pt(r_surf,nearest,eye,eye_mag) ! find nearest
! point on surface to the observer

c
c   Calculate the photosurface - store results in ph_surface
c   Calculate elevation angles at retarded position - store
c   values in sigh_tot.
c
call find_sigh_total(r_surf,ph_surface,sigh_tot,r_ret_hat,
1 nearest,eye,v,beta)

c
c   Finally do the distortion on the original surface
c   relative rotates each point on the cube as if it was a
c   microscopic cube.
c
call relative(surface,nearest,sigh_tot,r_surf,rotme,
1 r_ret_hat,eye,v,beta,pos_new,angle)

c
c   Relativistic rotation based upon elevation angle of the
c   retarded position
c
do isur = 1,6 ! new position of distorted surface
do ipnt = 1,16
do i = 1,3
  surface(isur,ipnt,i) = surface(isur,ipnt,i)
  1 + pos_new(i) ! move em out
enddo
enddo
enddo

c
c   Now finally,finally calculate Doppler shifted color
c   and map values into Dore (software color map) which is:
c   hue = 0.25 is green, hue = 0.50 is red, hue = 0.0 is blue
c   visible wavelengths:
c   λB(blue) = 390 nm, λG(green) = 500nm, λR(red) = 640 nm
c   Doppler: λ = λ0γ(1-βcosΨ) , let hue = ln(λ/λB)
c

```



```

c
a=0.
do i=1,3
  delta_r(i) = r_surf(is,ipnt,i) - nearest(i)
  a = delta_r(i)*nearest(i) + a
enddo
if ((a.lt.0.0).and.(a .lt. -0.1)) then
c Compute the square of the magnitude of delta_r
  bmag = 0.
  do i=1,3
    bmag=delta_r(i)*delta_r(i) + bmag
  enddo
c compute new nearest
  do i=1,3
    nearest(i)=nearest(i)-a*delta_r(i)/bmag
  enddo
endif
enddo

c
magn = 0.0
do i=1,3
  magn = nearest(i)*nearest(i) + magn
enddo
magn = sqrt(magn)
nearest(4)=magn      ! magnitude of nearest vector
100 return
end

subroutine rotate(surface,v,beta)
c
c Given cube at origin. Move to position(3) and then
c rotate it to point along the beta direction
c returns rotated array in surface
c
real beta,v(3)      ! v/c, velocity vector of cube
real surface(6,16,3) ! original/rotated array
real theta
real little        ! blowup protection
real A1,B1,C1      ! vector components of curl
real V1V,L         ! in transform matrices
real x1,x2,x3      ! temporary x position
real y1,y2,y3      ! temporary y position
real z1,z2,z3      ! temporary z position
real sinl,cosl,sinJ,cosJ ! angles for rotation about VXrip
real vector(3)     ! vector position of each pt on cube

c
little = 1.0e-6      ! to prevent blowup

c
c Determine axis of rotation.
c
A1 = 0.              ! components of rotation axis
B1 = v(3)/(beta + little)
C1 = -v(2)/(beta + little)

c
V1V = (B1*B1 + C1*C1)
L = sqrt(V1V + A1*A1)
V1V = sqrt(V1V)
cosl = C1/(V1V + little)
sinl = B1/(V1V + little)
cosJ = V1V/(L + little)
sinJ = A1/(L + little)

c
if((C1.eq.0.).and.(B1.eq.0.)) goto 100
c Rotation angle is found by taking dot product between y-axis
c and v.
c
theta = acos(v(1)/(beta + little))

c
do isur = 1,6
do ipnt = 1,16
do iu = 1,3
  vector(iu) = surface(isur,ipnt,iu)
enddo

c
x1 = vector(1)
y1 = cosl*vector(2)-sinl*vector(3)
z1 = cosl*vector(3)+sinl*vector(2)

c
c now rotate cube about y so that z axis corresponds to axis of
c rotation

```

```

c
x2 = cosJ*x1-sinJ*z1
y2 = y1
z2 = cosJ*z1+sinJ*x1

c
c now DO the relativistic rotation ( about new z)
c
x3 = x2*cos(theta) + y2*sin(theta)
y3 = y2*cos(theta) - x2*sin(theta)
z3 = z2

c
c now do inverse transforms
c
x2 = cosJ*x3+sinJ*z3      ! inverse rot about y
y2 = y3
z2 = cosJ*z3-sinJ*x3

c
x1 = x2                    ! inverse rot about x
y1 = cosl*y2+sinl*z2
z1 = cosl*z2-sinl*y2

c
c Now translate cube back to where it was at the outset
c
surface(isur,ipnt,1) = x1
surface(isur,ipnt,2) = y1
surface(isur,ipnt,3) = z1

c
enddo      ! ipnt over points
enddo      ! isur 6 surfaces

c
100 return
end

c
subroutine calc_r(Lor_surf,r_surf,beta,eye)
c
c Calculates the angle of rotation
c for each point on surface according to apparant relativistic
c rotation. Taylor Introductory Mechanics pg 357
c These are put into an array (sigh(6,16)) to be used later
c
real beta

```

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```

real eye(3)          ! position of observer's eye
real Lor_surf(6,16,3) ! points on surface of rotated cube
real little          ! overflow prevention
real r_surf(6,16,4)  ! for calculations (rpoint,rmag)
real rpoint(3),mag_rpoint

c
if(beta.eq.0)goto 100      ! bail out no need to work
little = 1.0e-6           ! divide protect

c
do is = 1,6                ! over all 6 surfaces
do ipnt = 1,16            ! 16 points/surface

c
Find the spatial location of points on the surface
as measured from the observer's position.

c
mag_rpoint = 0.0         ! initialize
do i = 1,3
  rpoint(i) = Lor_surf(is,ipnt,i) - eye(i)
  mag_rpoint = mag_rpoint + rpoint(i)**2
enddo
mag_rpoint = sqrt(mag_rpoint)

c
Calculate angle of elevation for this x,y,z triplet

c
do i = 1,3
  r_surf(is,ipnt,i) = rpoint(i)          ! save for later
enddo
r_surf(is,ipnt,4) = mag_rpoint
enddo                                     ! ipnt
enddo                                     ! is
100 return
end

subroutine find_sigh_total(r_surf,ph_surf,sigh_tot,r_ret_hat,
1 nearest,eye,v,beta)
c
real r_surf(6,16,4)      ! store scratch rpoint,magr,sigh
real ph_surf(6,16,3)     ! light surface - position of
c                          ! retarded emission points
real r_ret_hat(6,16,3)   ! unit vector along r-retarded
real beta,beta2,v(3),eye(3)
real rmag                ! magnitude of rpoint()
real nearest(4)          ! nearest point on cube to eye
real sigh_tot(6,16)     ! total rotation angle
real c,d,little
real a1,b1,c1,d1        ! for calculating rret
real r_dot_beta,b2rn,rmag,rret
real vt,ret_pos(3),cos_sigh

c
calculate angle beta makes with x-z plane

c
little = 1.0e-6         ! divide protect
if(beta.eq.0)goto 100   ! bail out no need to work
beta2 = beta*beta
rmag = nearest(4)
b2rn = beta2*rmag

c
do is = 1,6              ! over all 6 surfaces
do ipnt = 1,16          ! 16 points/surface
rmag = r_surf(is,ipnt,4)

c
Calculate the position of the points such that
emitted rays reach the observer simultaneous
with a ray from the nearest point.

c
r_dot_beta = 0.0
do i = 1,3
  r_dot_beta = r_surf(is,ipnt,i)*v(i) + r_dot_beta
enddo
a1 = 1. - beta2
b1 = 2*(r_dot_beta + b2rn)
c1 = -rmag*rmag - 2.*r_dot_beta*rmag-b2rn*rmag
d1 = b1*b1 - 4.*a1*c1
rret = (-b1 + sqrt(d1))/(2.*a1 + little)
vt = rret-rmag
cos_sigh = 0.0
do i = 1,3
  ret_pos(i) = r_surf(is,ipnt,i) -vt*v(i)
  r_ret_hat(is,ipnt,i) = ret_pos(i)/(rret + little)
  ph_surf(is,ipnt,i) = eye(i) + ret_pos(i)

```

```

c
Calculate the angle of elevation of the retarded position.
c
cos_sigh = - ret_pos(i)*v(i) + cos_sigh
enddo
cos_sigh = cos_sigh/(rret*beta + little)
sigh_tot(is,ipnt) = acos(cos_sigh)

c
enddo                                     ! ipnt
enddo                                     ! is

c
100 return
end

```

```

subroutine relative(surface,nearest,sigh_tot,r_surf,rotme,
1 r_ret_hat,eye,v,beta,pos_new,angle)
c
Rotate each point on surface according to apparant relativistic
rotation. Taylor + plane wave correction
c
With sigh_tot the total rotation angle move cube so that axis
through nearest point is z axis and rotate each point by sigh_tot
about this. rotate pointson the uncontracted cube. Points on the
contracted cube were used to find the rotation angles

c
real r_surf(6,16,4)      ! store scratch rpoint,magr,sigh
real beta,v(3),beta2,gamma ! v/c,velocity, beta*beta
real surface(6,16,3)     ! points on surface of rotated cube
real nearest(4)          ! nearest point on cube to eye
real sigh_tot(6,16)     ! total rotation angle
real rotme(6,16,3)      ! rotation axis
real r_ret_hat(6,16,3)   ! unit vector along r-retarded
real little              ! small number
real eye(3),pos_new(3)   ! dist eye to (o,o,o) and new pos
real ctr_to_nearest(3)   ! vector for moving cube b/4 rot
real ctn_unc(3)          ! uncontracted ctr_to_nearest
real proj(3),pro        ! projection along beta
real A1,B1,C1            ! vector components of curl
real V1V,L               ! in transform matrices
real theta,xx,angle      ! cos(total angle of elevation)
real ctn_dotv            ! cube center dot v()
real cos_doppler         ! cosine of doppler angle
real sign,signdop        ! sign's of angles
real sinI,cosI,sinJ,cosJ ! angles for rotation about VXrip
real vector(3)           ! vector pos of each pt on cube
real r_dot_rhat          ! dot product of r and r_ret_hat
real no_points           ! no points on cube

c
if(beta.eq.0.0)goto 100 ! no action dont bother
little = 1.0e-6
no_points = 16.0*6.0
beta2 = beta*beta
gamma = 1./sqrt(1.-beta2))
ctn_dotv = 0.          ! used to undo a Lorentz contract
do i = 1,3
  ctr_to_nearest(i) = nearest(i) - pos_new(i) + eye(i)
  ctn_dotv = ctr_to_nearest(i)*v(i) + ctn_dotv
enddo
ctn_dotv = ctn_dotv*(gamma - 1.0)/beta2

c
Uncontract center-to-nearest point on cube vector

c
do i = 1,3
  ctn_unc(i) = ctr_to_nearest(i) + ctn_dotv*v(i)
enddo

c
Rotate each vector on the surface around rotme by angle sigh_tot
do is = 1,6
do ipnt = 1,16
  A1 = rotme(is,ipnt,1)
  B1 = rotme(is,ipnt,2)
  C1 = rotme(is,ipnt,3)
  V1V = (B1*B1 + C1*C1)
  L = sqrt(V1V + A1*A1)
  V1V = sqrt(V1V)
  cosI = C1/(V1V + little)
  sinI = B1/(V1V + little)
  cosJ = V1V/(L + little)
  sinJ = A1/(L + little)

c
Calculate angle of elevation for this x,y,z triplet
For Doppler shift evaluation

```

```

c      cos_doppler = 0.0
      do i = 1,3
        cos_doppler = cos_doppler + r_surf(is,ipnt,i)*v(i)
      enddo
      cos_doppler = cos_doppler/(r_surf(is,ipnt,4) + little)
      angle = angle + cos_doppler          ! ave angle for
                                          ! Doppler shift
c
c      Now do relativistic rotation
c
c      sigh = sigh_tot(is,ipnt)
      cos_sigh = cos(sigh)
      xx = (cos_sigh - beta)/(1. - beta*cos_sigh + little)
      if(abs(xx).gt.1.) xx = 1.
      theta = acos(xx) - sigh
c
c      Before rotation about rotme we must
c      Translate cube back so that it will rotate about nearest()
c      Recalculate Lorentz contraction
c      Uncontract component of ctr_to_nearest along velocity vector
      do i = 1,3
c      Vector points from a point on the surface nearest the
c      observer to each labeled point on the surface
        vector(i) = surface(is,ipnt,i) - ctr_unc(i)
      enddo
c
c      now rotate cube (about x) so that new axis of rotation
c      is in the x-z plane
c
c      x1 = vector(1)
      y1 = cosI*vector(2)-sinI*vector(3)
      z1 = cosI*vector(3)+sinI*vector(2)
c
c      rotate cube about y so that z axis corresponds to axis of rotation
c
      x2 = cosJ*x1-sinJ*z1
      y2 = y1
      z2 = cosJ*z1+sinJ*x1
c
c      now DO the relativistic rotation ( about new z)
c
c      x3 = x2*cos(theta) + y2*sin(theta)
      y3 = y2*cos(theta) - x2*sin(theta)
      z3 = z2
c
c      now do inverse transforms
c
c      x2 = cosJ*x3+sinJ*z3          ! inverse rot about y
      y2 = y3
      z2 = cosJ*z3-sinJ*x3
c
c      x1 = x2          ! inverse rot about x
      y1 = cosI*y2+sinI*z2
      z1 = cosI*z2-sinI*y2
c
c      Correct rotated surface so that each point lies along direction
c      of corresponding point on photosurface.
c
      r_dot_rhat = (nearest(1) + x1)*r_ret_hat(is,ipnt,1)
      r_dot_rhat = (nearest(2) + y1)*r_ret_hat(is,ipnt,2)
      l          + r_dot_rhat
      r_dot_rhat = (nearest(3) + z1)*r_ret_hat(is,ipnt,3)
      l          + r_dot_rhat
      x1 = r_dot_rhat*r_ret_hat(is,ipnt,1) - nearest(1)
      y1 = r_dot_rhat*r_ret_hat(is,ipnt,2) - nearest(2)
      z1 = r_dot_rhat*r_ret_hat(is,ipnt,3) - nearest(3)
c
c      Now translate cube back to where it was at the outset
      surface(is,ipnt,1) = x1 + ctr_to_nearest(1)
      surface(is,ipnt,2) = y1 + ctr_to_nearest(2)
      surface(is,ipnt,3) = z1 + ctr_to_nearest(3)
c
c      enddo          ! ipnt
      enddo          ! is
      angle = angle/no_points ! ave angle cube makes with x-z plane
      return
      end
100

```

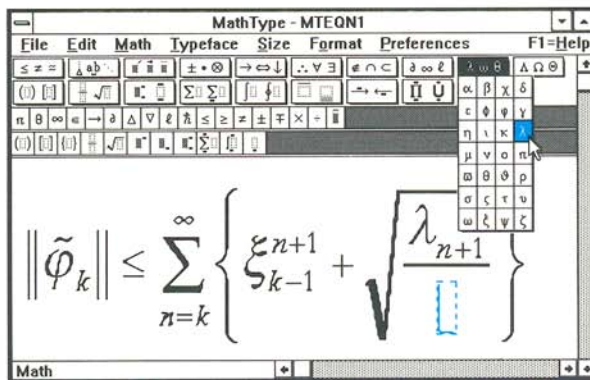
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